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## **New Implementation of Reproducing Kernel Hilbert Space Method for Solving a Class of Third-order Differential Equations**

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### ***Abstract***

In this paper, we apply the new implementation of reproducing kernel Hilbert space method to give the approximate solution to some third-order boundaryvalue problems with variable coefficients. In this method, the analytical solution is expressed in the form of a series. At the end, two examples are given to illustrate implementation, accuracy and effectiveness of the method.

**Keywords:** Reproducing kernel Hilbert space method; Boundary value problems; Third-order differential equations; Approximatesolution.

### **1. Introduction**

Reproducing kernel Hilbert space method is a promising method which has been applied more and more for solving various problems such as ordinary differential equations, partial differential equations, differential-difference equations, integral equations, and etc. in the previous decades [1]-[22]. Approximate solution of the Fredholm integral equation of the first kind in the

reproducing kernel Hilbert space was presented by Du and Cui [3,4], solution of a system of linear Volterra integral equations was discussed by Yang et al. [5], solvability of a class of Volterra integral equations with weakly singular kernel using reproducing kernel Hilbert space method were investigated in [6,7,8], Geng [9] explained how to solve the Fredholm integral equation of the third kind in the reproducing kernel Hilbert space method. These are a bunch of extensive works related to reproducing kernel Hilbert space method for solving integral equations.

In 1986, Cui Minggen [10] gave the reproducing kernel space  $W_2^1[a, b]$  and its reproducing kernel. This technique has successfully been treated singular linear two-point boundary value problems [11,12], singular nonlinear second-order boundary value problems [13,14,15,16], nonlinear system of boundary value problems [17], third-order boundary value problems [18,19], fifth-order boundary value problems [20], and nonlinear partial differential equations [21] in recent years.

This paper investigates the approximate solution of the following third-order boundary value problem using new implementation of reproducing kernel Hilbert space method

$$\begin{cases} y'''(x) + p(x)y(x) = f(x), & 0 \leq x \leq 1, \\ y(0) = A, \quad y'(0) = B, \quad y'(1) = C, \end{cases} \quad (1)$$

where  $p(x)$ ,  $f(x)$  are analytical known functions defined on the interval  $[0,1]$ , unknown function  $y(x)$  is continuous on the interval  $[0,1]$  and  $A, B, C$  are finite real constants.

Several numerical techniques have been proposed to solve high-order differential equations [23,24,25].

As we know, Gram-Schmidt orthogonalization process is numerically unstable and in addition it may take a lot of time to produce numerical approximation. Here, instead of using orthogonal process, we successfully make use of the basic functions which are obtained by reproducing kernel Hilbert space method.

This paper is organized as follows. In the following section, we introduce some useful definitions and theorems. Section 3 is devoted to solve Eq. (1) by new implementation of reproducing kernel Hilbert space method. Two numerical examples are presented in Section 4. We end the paper with a few conclusions.

## 2. Reproducing Kernel Spaces

In this section, we follow the recent work of [1]-[22] and represent some useful materials.

### Definition 1.

Let  $(\mathcal{H}, \langle \cdot, \cdot \rangle_{\mathcal{H}})$  be a Hilbert space of real-valued functions on some nonempty set  $\mathcal{X}$ . A function  $k: \mathcal{X} \times \mathcal{X} \rightarrow \mathbb{R}$  is said to be the *reproducing kernel* of  $\mathcal{H}$  if and only if

1.  $k(x, \cdot) \in \mathcal{H}, \quad \forall x \in \mathcal{X},$
2.  $\langle \varphi(\cdot), k(x, \cdot) \rangle_{\mathcal{H}} = \varphi(x), \quad \forall \varphi \in \mathcal{H}, \quad \forall x \in \mathcal{X}, \quad (\text{reproducing property}).$

It is known that the reproducing kernel of a reproducing kernel Hilbertspace is unique and the existence of a reproducing kernel is according to the *Riesz Representation Theorem*. The reproducing kernel  $k$  of a Hilbert space  $\mathcal{H}$  quite determines the space  $\mathcal{H}$ . Each set of functions  $\{\varphi_i\}_{i=1}^\infty$  which converges strongly to a function  $\varphi$  in  $\mathcal{H}$ , converges also in the pointwise sense. In addition, this convergence is uniform on every subset of  $\mathcal{X}$  on which  $x \mapsto k(x, x)$  is bounded.

**Definition 2.**

$$W_2^4[0,1] = \{y(x) | y'''(x) \text{ is an absolute continuous real valued function and } y^{(4)}(x) \in L^2[0,1], y(0) = y'(0) = y'(1) = 0\}.$$

The inner product and the norm in the function space  $W_2^4[0,1]$  are defined as follows.

$$\begin{cases} \langle u, v \rangle_{W_2^4} = u''(0)v''(0) + \int_0^1 u^{(4)}(x)v^{(4)}(x)dx, \\ \|u\|_{W_2^4} = \sqrt{\langle u, u \rangle_{W_2^4}}. \end{cases}$$

Let's assume that function  $R(x, t) \in W_2^4[0,1]$  satisfies the following generalized differential equations

$$\begin{cases} \frac{\partial^8 R(x, t)}{\partial t^8} = \delta(t - x), \\ R(x, 1) - \frac{\partial^7 R(x, 1)}{\partial t^7} = 0, \\ \frac{\partial^4 R(x, 1)}{\partial t^4} = 0, \quad \frac{\partial^4 R(x, 0)}{\partial t^4}, \\ \frac{\partial^5 R(x, 1)}{\partial t^5} = 0, \quad \frac{\partial^5 R(x, 0)}{\partial t^5}. \end{cases} \tag{2}$$

where  $\delta$  is the Dirac delta function. Therefore, the following theorem holds.

**Theorem 1.** Under the assumptions of Eq. (2), Hilbert space  $W_2^4[0,1]$  is a reproducing kernel Hilbertspace with the reproducing kernel function  $R(x, t)$ , namely for any  $y(t) \in W_2^4[0,1]$  and each fixed  $x \in [0,1]$ , there exists  $R(x, t) \in W_2^4[0,1]$ ,  $t \in [0,1]$ , such that

$$\langle y(\cdot), R(x, \cdot) \rangle_{W_2^4} = y(x).$$

While  $x \neq t$ , function  $R(x, t)$  is the solution of the following constant linear homogeneous differential equation with 8 orders,

$$\frac{\partial^8 R(x, t)}{\partial t^8} = 0, \tag{3}$$

with the boundary conditions:

$$\begin{cases} R(x, 1) - \frac{\partial^7 R(x, 1)}{\partial t^7} = 0, \\ \frac{\partial^4 R(x, 1)}{\partial t^4} = 0, \quad \frac{\partial^4 R(x, 0)}{\partial t^4}, \\ \frac{\partial^5 R(x, 1)}{\partial t^5} = 0, \quad \frac{\partial^5 R(x, 0)}{\partial t^5}. \end{cases} \quad (4)$$

We know that Eq. (3) has characteristic equation  $\lambda^8 = 0$ , and the eigenvalue  $\lambda = 0$  is a root whose multiplicity is 8. Hence, the general solution of Eq. (2) is

$$R(x, t) = \begin{cases} \sum_{i=1}^8 c_i(x) t^{i-1}, & t \leq x, \\ \sum_{i=1}^8 d_i(x) t^{i-1}, & t > x. \end{cases} \quad (5)$$

Now, we are ready to calculate the coefficients  $c_i(x)$  and  $d_i(x)$ ,  $i = 1, \dots, 8$ . Since

$$\frac{\partial^8 R(x, t)}{\partial t^8} = \delta(t - x),$$

we have

$$\begin{cases} \frac{\partial^k R(x, x^+)}{\partial t^k} = \frac{\partial^k R(x, x^-)}{\partial t^k}, & k = 0, \dots, 6, \\ \frac{\partial^7 R(x, x^+)}{\partial t^7} - \frac{\partial^7 R(x, x^-)}{\partial t^7} = 1. \end{cases} \quad (6)$$

Then, using Eqs. (4) and (6), the unknown coefficients of Eq. (5) are uniquely obtained (in Appendix A).

**Definition 3.**

$$W_2^1[0,1]$$

= { $y(x)$  |  $y(x)$  is an absolute continuous real valued function on the interval  $[0,1]$  and  $y'(x) \in L^2[0,1]$ }.

The inner product and the norm in the function space  $W_2^1[0,1]$  are defined as follows.

$$\begin{cases} \langle u, v \rangle_{W_2^1} = u(0)v(0) + \int_0^1 u'(x)v'(x)dx, \\ \|u\|_{W_2^1} = \sqrt{\langle u, u \rangle_{W_2^1}}. \end{cases}$$

**Theorem 2.** Hilbert space  $W_2^1[0,1]$  is a reproducing kernel space with the reproducing kernel function

$$Q(x, t) = \begin{cases} 1 + t, & t \leq x, \\ 1 + x, & t > x, \end{cases}$$

that is, for any  $y(t) \in W_2^1[0,1]$  and each fixed  $x \in [0,1]$ , it follows that

$$\langle y(\cdot), Q(x, \cdot) \rangle_{W_2^1} = y(x).$$

### 3. Reproducing Kernel Hilbert Space Method

We suppose that Eq. (1) has a unique solution. To deal with the system, we consider Eq. (1) as

$$\mathbb{L}y(x) = f(x), \quad 0 \leq x \leq 1, \tag{7}$$

where  $\mathbb{L}y(x) = y'''(x) + p(x)y(x)$ , it is clear that  $\mathbb{L}$  is the bounded linear operator of  $W_2^4[0,1] \rightarrow W_2^1[0,1]$ . We shall give the representation of analytical solution of Eq. (7) in the space  $W_2^4[0,1]$ . Set  $\varphi_i(x) = Q(x_i, x)$  and  $\psi_i(x) = \mathbb{L}^*\varphi_i(x)$ ,  $i = 1, 2, \dots$ , where  $Q(x_i, x)$  is the reproducing kernel of  $W_2^1[0,1]$  and  $\mathbb{L}^*$  is the adjoint operator of  $\mathbb{L}$ .

**Theorem 3.** Let  $\{x_i\}_{i=1}^\infty$  be a dense subset of interval  $[0,1]$ , then  $\{\psi_i(x)\}_{i=1}^\infty$  is a complete system of  $W_2^4[0,1]$  and  $\psi_i(x) = \mathbb{L}_t R(x, t)|_{t=x_i}$ , where the subscript  $t$  in the operator  $\mathbb{L}$  indicates that the operator  $\mathbb{L}$  applies to the function of  $t$ .

Usually, a normalized orthogonal system is constructed from  $\{\psi_i(x)\}_{i=1}^\infty$  by using the Gram-Schmidt algorithm, and then the approximate solution be obtained by calculating a truncated series based on these functions. However, Gram-Schmidt algorithm has some drawbacks such as numerical instability and high volume of computations. Here, to fix these flaws, we state the following Theorem in which the following notation are used.

$$\hat{\mathbf{a}} = \begin{bmatrix} \hat{a}_1 \\ \hat{a}_2 \\ \vdots \\ \hat{a}_N \end{bmatrix}, \quad \bar{\mathbf{a}} = \begin{bmatrix} \bar{a}_1 \\ \bar{a}_2 \\ \vdots \\ \bar{a}_N \end{bmatrix}, \quad \bar{\mathbf{F}} = \begin{bmatrix} \hat{f}_1 \\ \hat{f}_2 \\ \vdots \\ \hat{f}_N \end{bmatrix}, \quad B = \begin{bmatrix} \beta_{11} & 0 & \dots & 0 \\ \beta_{21} & \beta_{22} & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ \beta_{N1} & \beta_{N2} & \dots & \beta_{NN} \end{bmatrix},$$

$$\Psi = \begin{bmatrix} \psi_{11} & \psi_{12} & \dots & \psi_{1N} \\ \psi_{21} & \psi_{22} & & \psi_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ \psi_{N1} & \psi_{N2} & \dots & \psi_{NN} \end{bmatrix}, \quad \tilde{\Psi} = \begin{bmatrix} \tilde{\psi}_{11} & \tilde{\psi}_{12} & \dots & \tilde{\psi}_{1N} \\ \tilde{\psi}_{21} & \tilde{\psi}_{22} & & \tilde{\psi}_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ \tilde{\psi}_{N1} & \tilde{\psi}_{N2} & \dots & \tilde{\psi}_{NN} \end{bmatrix},$$

where  $\psi_{ij} = \langle \mathbb{L}\psi_j, \psi_i \rangle$ ,  $\tilde{\psi}_{ij} = \langle \mathbb{L}\psi_j, \tilde{\psi}_i \rangle$ ,  $\hat{f}_i = \langle f, \psi_i \rangle$ ,  $\beta_{ii} > 0$ ,  $i, j = 1, \dots, N$ .

**Theorem 4.** Suppose that  $\{\psi_i(x)\}_{i=1}^\infty$  a linearly independent set in  $W_2^4[0,1]$  and  $\{\tilde{\psi}_i(x)\}_{i=1}^\infty$  be a normalized orthogonal system in  $W_2^4[0,1]$ , such that  $\tilde{\psi}_i(x) = \sum_{k=1}^i \beta_{ik} \psi_k(x)$ . If  $y(x) = \sum_{i=1}^\infty \bar{a}_i \psi_i(x) \simeq y_N(x) = \sum_{i=1}^N \bar{a}_i \psi_i(x) = \sum_{i=1}^N \hat{a}_i \tilde{\psi}_i(x)$  then  $\Psi \bar{\mathbf{a}} = \bar{\mathbf{F}}$ .

**Proof.** Suppose that  $y(x) \in W_2^4[0,1]$  then  $y(x) = \sum_{i=1}^\infty \bar{a}_i \psi_i(x) = \sum_{i=1}^\infty \hat{a}_i \tilde{\psi}_i(x)$ . Now, by truncating N-term of the two series, because of  $y_N(x) = \sum_{i=1}^N \bar{a}_i \psi_i(x) = \sum_{i=1}^N \hat{a}_i \tilde{\psi}_i(x)$  and since  $\tilde{\psi}_i(x) = \sum_{k=1}^i \beta_{ik} \psi_k(x)$  so

$$\sum_{i=1}^N \bar{a}_i \psi_i(x) = \sum_{i=1}^N \hat{a}_i \tilde{\psi}_i(x) = \sum_{i=1}^N \hat{a}_i \left( \sum_{k=1}^i \beta_{ik} \psi_k(x) \right) = \sum_{k=1}^N \left( \sum_{i=k}^N \hat{a}_i \beta_{ik} \right) \psi_k(x).$$

Due to the linear independence of  $\{\psi_i(x)\}_{i=1}^\infty, \bar{a}_k = \sum_{i=k}^N \hat{a}_i \beta_{ik}, k = 1, \dots, N$  therefore

$$\bar{\mathbf{a}} = B^T \hat{\mathbf{a}}. \tag{8}$$

Eq. (7), imply  $\mathbb{L}y_N(x) = f(x)$ . For  $i = 1, \dots, N$  we have

$$\begin{aligned} \langle \mathbb{L}y_N, \tilde{\psi}_i \rangle &= \langle f, \tilde{\psi}_i \rangle \Rightarrow \sum_{j=1}^N \hat{a}_j \langle \mathbb{L}\tilde{\psi}_j, \tilde{\psi}_i \rangle = \langle f, \tilde{\psi}_i \rangle \\ &\Rightarrow \sum_{j=1}^N \hat{a}_j \sum_{k=1}^i \beta_{ik} \sum_{l=1}^j \beta_{jl} \langle \mathbb{L}\psi_l, \psi_k \rangle = \sum_{k=1}^i \beta_{ik} \langle f, \psi_k \rangle \\ &\Rightarrow \sum_{j=1}^N \hat{a}_j \sum_{k=1}^i \sum_{l=1}^j \beta_{ik} \langle \mathbb{L}\psi_l, \psi_k \rangle \beta_{lj}^T = \sum_{k=1}^i \beta_{ik} \langle f, \psi_k \rangle \\ &\Rightarrow \sum_{j=1}^N \hat{a}_j (B \Psi B^T)_{ij} = \sum_{k=1}^i \beta_{ik} \langle f, \psi_k \rangle \\ &\Rightarrow (B \Psi B^T) \hat{\mathbf{a}} = B \bar{\mathbf{F}}. \end{aligned}$$

Eq. (8), imply  $B \Psi \bar{\mathbf{a}} = B \bar{\mathbf{F}}$ , hence

$$\Psi \bar{\mathbf{a}} = \bar{\mathbf{F}}. \quad \blacksquare$$

It is necessary to mention that here we solve the system  $\Psi \bar{\mathbf{a}} = \bar{\mathbf{F}}$  which obtained without using the Gram-Schmidt algorithm.

#### 4. Numerical examples

To illustrate the effectiveness and the accuracy of the proposed method, two numerical examples are considered in this section. Figures 1,2,3 and 4 show that the approximate solution and its

derivatives up to third-order, converge to the exact solution and its derivatives. We solved these examples by the reproducing kernel Hilbert space method with  $x_i = \frac{i-1}{N-1}, i = 1, \dots, N$  for  $N = 5$ .

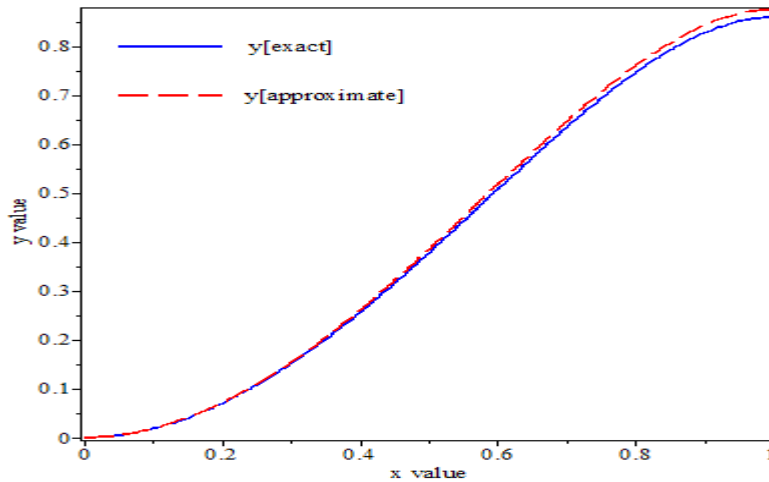


Figure 1: Comparison between approximate solution and the exact solution for Example 1 for  $N = 5$ .

**Example 1.** We consider the following third-order BVP

$$\begin{cases} y'''(x) - xy(x) = (x^3 - 2x^2 - 5x - 3)e^x, & 0 \leq x \leq 1, \\ y(0) = 0, \quad y'(0) = 1, \quad y'(1) = -e. \end{cases}$$

The exact solution of the above system is  $y(x) = x(1 - x)e^x$ .

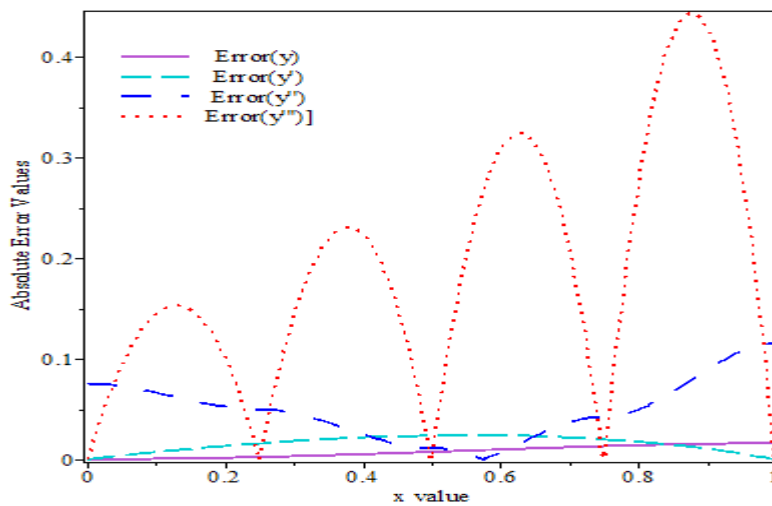


Figure 2: The absolute error between approximate solution and the exact solution and its derivatives for Example 1 for  $N = 5$ .

**Example 2.** Consider the following third-order BVP

$$\begin{cases} y'''(x) - xy(x) = (1 - x)e^x, & 0 \leq x \leq 1, \\ y(0) = 1, \quad y'(0) = 1, \quad y'(1) = e. \end{cases}$$

The exact solution of the above system is  $y(x) = e^x$ .

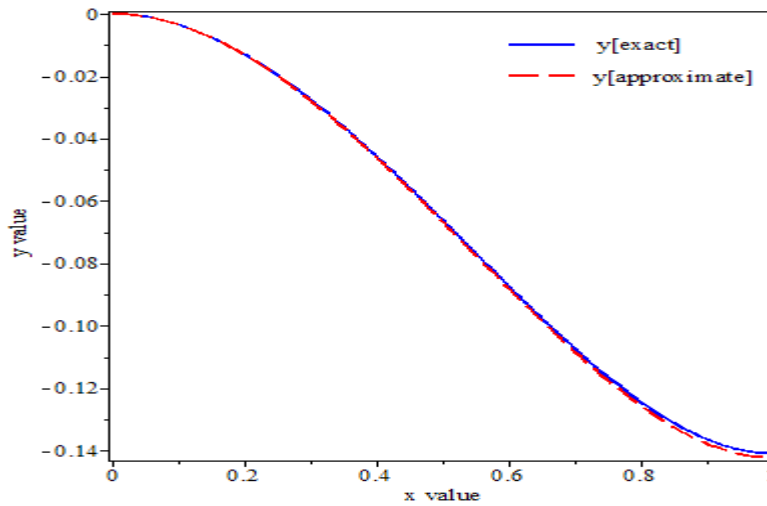


Figure 3: Comparison between approximate solution and the exact solution for Example 2 for  $N = 5$ .

#### 4. Conclusions

In this paper, we introduced the new implementation of reproducing kernel Hilbert space method to obtain the approximate solution to some third-order boundary value problems with variable coefficients. The reliability of the method and reduction of the amount of computation gives this method a wider applicability. The obtained numerical results confirm that the method is rapidly convergent and show that the approximate solution converge to the exact solution.

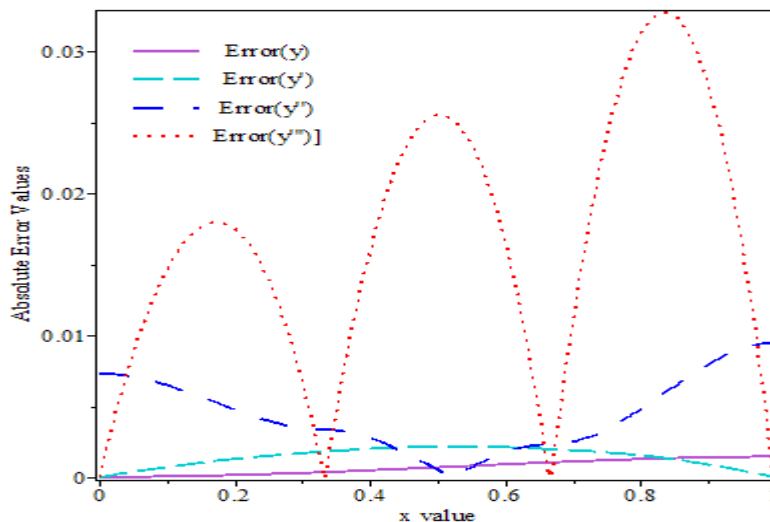


Figure 4: The absolute error between approximate solution and the exact solution and its derivatives for Example 2 for  $N = 5$ .



**Appendix A.**

$$\begin{aligned}
c_1 = c_2 = 0, \quad c_3 &= \frac{x^2}{7!}(-14x^4 + 3x^5 + 45372 - 30262x + 21x^3), \quad c_4 \\
&= -\frac{x^2}{7!}(2x^5 + 30262 - 20212x - 7x^4 + 35x^2), \\
c_5 = c_6 = 0, \quad c_7 &= \frac{x(x-1)^2}{6!}, \quad c_8 = -\frac{(2x+1)(x-1)^2}{7!}, \quad d_1 = -\frac{x^7}{7!}, \quad d_2 = \frac{x^6}{6!}, \quad d_3 \\
&= \frac{x^2}{7!}(-14x^4 + 3x^5 + 45372 - 30262x), \quad d_4 \\
&= -\frac{x^2}{7!}(2x^5 + 30262 - 20212x - 7x^4), \quad d_5 = -\frac{x^3}{3!4!}, \quad d_6 = \frac{x^2}{2!5!}, \quad d_7 \\
&= \frac{x^2(x-2)}{6!}, \quad d_8 = \frac{x^2(3-2x)}{7!}.
\end{aligned}$$

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