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# Some integrability estimates for solutions of the fractional p-Laplace equation

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### Abstract

For  $(\alpha, p) \in (0, 1) \times (1, \infty)$ , this note focuses on some integrability estimates for solutions of the following Dirichlet problem

$$\begin{cases} \mathsf{L}_{\alpha,p}\mathfrak{u}(x)=\mathfrak{g}(x) & \text{ as } x\in\Omega,\\ \mathfrak{u}(x)=0 & \text{ as } x\in\mathbb{R}^n\backslash\Omega, \end{cases}$$

where  $L_{\alpha,p}$  is the fractional p-Laplace operator. ©2017 All rights reserved.

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# 1. Introduction

Unless stated otherwise, we will always assume that  $(\alpha, p) \in (0, 1) \times (1, \infty)$  and  $\Omega$  is a bounded Lipschitz domain. This paper is devoted to a further study of the integrability estimates for weak solutions of the following Dirichlet problem

$$\begin{cases} L_{\alpha,p}u(x) = g(x) & \text{ as } x \in \Omega, \\ u(x) = 0 & \text{ as } x \in \mathbb{R}^n \setminus \Omega. \end{cases}$$
(1.1)

Here  $L_{\alpha,p}$  is the so-called fractional p-Laplace operator

$$L_{\alpha,p}\mathfrak{u}(x)=p.v.\int_{\mathbb{R}^n}\frac{|\mathfrak{u}(y)-\mathfrak{u}(x)|^{p-2}(\mathfrak{u}(y)-\mathfrak{u}(x))}{|y-x|^{n+\alpha p}}dy.$$

When p = 2,  $L_{\alpha,2}$  has already been known as the classical fractional Laplace operator, which has initially been studied ([11, 18]). It is a generator of a strongly continuous contractive semigroup on  $L^2(\mathbb{R}^n)$  that can be extended to contraction semigroup on  $L^p(\mathbb{R}^n)$  for  $p \in [1, \infty]$  ([3, 8]). The Dirichlet boundary problem of  $L_{\alpha,2}$  has been intensively investigated and many fundamental results have been proved, we refer the reader to [2, 4, 8, 12, 14] and the references therein for a fuller treatment of this topic. As a nonlinear generalization of  $L_{\alpha,2}$ ,  $L_{\alpha,p}$  has been extensively explored in recent years ([1, 5, 10]).

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When  $g = -\lambda |u|^{p-2}u$  with  $\lambda > 0$ , the equation

$$\mathcal{L}_{\alpha,p}\mathfrak{u}(\mathbf{x}) = -\lambda|\mathfrak{u}(\mathbf{x})|^{p-2}\mathfrak{u}(\mathbf{x}),\tag{1.2}$$

which also called the nonlocal Euler-Lagrange equation was fully discussed in [13] for large values of p and the limit equation as  $p \to \infty$  was derived. Equation (1.2) was closely related to the nonlocal eigenvalue problem and its viscosity solutions have many interesting properties. If  $g = -|u(x)|^{p-2}u(x)$ , a local version of (1.1), i.e.,

$$\begin{cases} L_{\alpha,p}u(x) = -|u(x)|^{p-2}u(x) & \text{ as } x \in \Omega, \\ u(x) = 1 & \text{ as } x \in K, \text{ compact } K \subset \Omega, \\ u(x) = 0 & \text{ as } x \in \mathbb{R}^n \backslash \Omega, \end{cases}$$
(1.3)

was studied in [15, 16]. It was proved that the weak solution of (1.3), which was nothing but the viscosity solution, was the capacitary potential of the relative fractional Sobolev capacity.

In [1], Barrios et al. studied the summability of the finite energy solutions to (1.1) in terms of the summability of g for  $\alpha p < n$  by adapting the ideas used in [12] for p = 2. In this paper, highly inspired by their methods and some known estimates for bounded Lipschitz domain, we obtain the following results for  $\alpha p \ge n$ .

**Theorem 1.1.** Let  $g \in L^{s}(\Omega)$ ,  $p \leq q < \infty$  with  $\alpha p = n$  and  $\Omega$  be a  $W^{\alpha,p}$ -extension domain. Then the solution of (1.1), denoted by u, satisfies the following boundedness:

(a) for 
$$s > \frac{\alpha q}{\alpha q - n}$$
, there exists a constant  $C_1 := C_1(n, \alpha, q, \Omega, \|g\|_{L^s(\Omega)})$  such that  $\|u\|_{L^{\infty}(\Omega)} \leq C_1$ ;

(b) for  $s = \frac{\alpha q}{\alpha q - n}$ , there exists a constant  $C_2 := C_2(n, \alpha, q, \Omega, \|g\|_{L^s(\Omega)})$  such that

$$\int_{\Omega} \exp^{\beta |u(x)|} dx \leqslant C_2, \quad \textit{for some} \quad \beta > 0;$$

(c) for  $\frac{q}{q-1} \leqslant s < \frac{\alpha q}{\alpha q-n}$ , there exists a constant  $C_3 := C_3(n, \alpha, s)$  such that

$$\|\mathbf{u}\|_{L^{s^*}(\Omega)} \leqslant C_3 \|g\|_{L^s(\Omega)}^{\frac{\alpha}{n-\alpha}} \quad as \quad s^* = \frac{qs(n-\alpha)}{s(n-q\alpha)+q\alpha}$$

**Theorem 1.2.** Assume that  $g \in L^1(\Omega)$ ,  $\alpha p > n$  and u is a solution of (1.1). Then there exists a constant  $C := C(n, \alpha, p, \Omega, ||g||_{L^1(\Omega)})$  such that

$$\|\mathfrak{u}\|_{L^{\infty}(\Omega)} \leq C.$$

We end this section with the outline of this paper. Section 2 presents some basic definitions and preliminaries. In Section 3, we give the proofs of Theorem 1.1 and Theorem 1.2.

# 2. Preliminaries

We recall that the inhomogeneous fractional Sobolev space  $W^{\alpha,p}(\mathbb{R}^n)$  is defined as

$$W^{\alpha,p}(\mathbb{R}^n) := \left\{ f \in L^p(\mathbb{R}^n) : \frac{|f(y) - f(x)|}{|y - x|^{\alpha + n/p}} \in L^p(\mathbb{R}^n \times \mathbb{R}^n) \right\},\$$

endowed with the norm

$$|f||_{W^{\alpha,p}(\mathbb{R}^n)} := \left( \int_{\mathbb{R}^n} |f(x)|^p dx + \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \frac{|f(y) - f(x)|^p}{|y - x|^{n + \alpha p}} dy dx \right)^{1/p}$$

Also,  $W^{\alpha,p}(\Omega)$  can be defined similarly with  $\mathbb{R}^n$  replaced by  $\Omega$ . The homogeneous fractional Sobolev space  $\dot{W}^{\alpha,p}(\mathbb{R}^n)$  can be defined by the semi-norm

$$\|f\|_{\dot{W}^{\alpha,p}(\mathbb{R}^n)} := \left( \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \frac{|f(y) - f(x)|^p}{|y - x|^{n + \alpha p}} dy dx \right)^{1/p}$$

which is the so-called Gagliardo semi-norm of f. Fractional Sobolev spaces have been a classical topic in functional and harmonic analysis all along, see e.g., the review paper [6] and the references therein.  $C(\Omega)$  is the space of all real-valued and continuous functions on  $\Omega$ . For each natural number k, i.e.,  $k \in \mathbb{N}$ ,  $C^k(\Omega)$  denotes the space of all functions being k times continuously differentiable,  $C_c^k(\Omega)$  stands for the space of all functions in  $C^k(\Omega)$  having compact support.  $C_c^{\infty}(\Omega)$  is the subspace of  $C_c^k(\Omega)$  given by  $C_c^{\infty}(\Omega) := \bigcap_k C_c^k(\Omega)$ .  $W_0^{\alpha,p}(\Omega)$  is the closure of  $C_c^{\infty}(\Omega)$  in the norm  $\|\cdot\|_{W^{\alpha,p}(\Omega)}$ .

Theorem 1.1 is based on the assumption that  $\Omega$  is an extension domain. We say that  $\Omega$  is a  $W^{\alpha,p}$ -extension domain, if there is a positive constant  $C := C(n, p, \Omega, \alpha)$  such that for every function  $f \in W^{\alpha,p}(\Omega)$ , there exists a function  $\tilde{f} \in W^{\alpha,p}(\mathbb{R}^n)$  with

$$f(x) = f(x)$$
 as  $x \in \Omega$ , and  $||f||_{W^{\alpha,p}(\mathbb{R}^n)} \leq C ||f||_{W^{\alpha,p}(\Omega)}$ 

Fractional extension results are essential to improve some fractional embedding theorems and have been discussed by many people such as Nezza-Palatulli-Valdinoci [6], Shvartsman [17], Triebel [19] and Zhou [21]. It is well-known that the space  $W^{\alpha,p}(\Omega)$  is continuously embedded in  $L^q(\Omega)$  for any  $q \in [p, \infty)$  when  $\Omega$  is a  $W^{\alpha,p}$ -extension domain. That is:

**Lemma 2.1.** Let  $\alpha p = n$ ,  $q \in [p, \infty)$  and  $\Omega$  be a  $W^{\alpha,p}$ -extension domain. Then there is a constant

$$C := C(\Omega, n, \alpha, p) > 0,$$

such that

$$\|f\|_{L^{q}(\Omega)} \leq C \|f\|_{W^{\alpha,p}(\Omega)}, \quad \forall f \in W^{\alpha,p}(\Omega)$$

*Moreover, if*  $f \in \dot{W}^{\alpha,p}(\Omega) \cap C_{c}(\Omega)$ *, then* 

$$\|f\|_{L^{q}(\Omega)} \leq C \|f\|_{\dot{W}^{\alpha,p}(\Omega)'}$$

where  $C_{c}(\Omega)$  stands for all real-valued and continuous functions on  $\Omega$  having compact support.

Proof. The first estimate is just [6, Theorem 6.10]. It only needs to prove the second one. If

$$f \in \dot{W}^{\alpha,p}(\Omega) \cap C_{c}(\Omega),$$

then supp f := K is a compact set, it follows from [9, Lemma 2.6] that  $f \in L^{p}(K)$ , i.e., there is a constant  $C := C(n, \alpha, p, \Omega) > 0$  such that

$$\|f\|_{L^p(\mathsf{K})} \leqslant C \|f\|_{\dot{W}^{\alpha,p}(\Omega)}.$$

Consequently, one has

$$\|f\|_{L^{q}(\Omega)} \leqslant C \|f\|_{W^{\alpha,p}(\Omega)} = C(\|f\|_{L^{p}(K)} + \|f\|_{\dot{W}^{\alpha,p}(\Omega)}) \leqslant C \|f\|_{\dot{W}^{\alpha,p}(\Omega)}.$$

A function  $u \in W^{\alpha,p}(\Omega)$  is a weak solution of (1.1) subject to the boundary condition u(x) = 0 on  $\mathbb{R}^n \setminus \Omega$  if

$$\langle \mathsf{L}_{\alpha,p}\mathfrak{u}, \phi \rangle = \langle \mathfrak{g}, \phi \rangle, \ \forall \phi \in W_0^{\alpha,p}(\Omega),$$

where

$$\langle \mathcal{L}_{\alpha,p}\mathfrak{u}, \phi \rangle := \int_{\Omega} \int_{\Omega} \frac{|\mathfrak{u}(y) - \mathfrak{u}(x)|^{p-2}(\mathfrak{u}(y) - \mathfrak{u}(x))(\phi(y) - \phi(x))}{|y - x|^{n+\alpha p}} dy dx,$$
(2.1)

and  $\langle g, \phi \rangle$  is given by the duality product. It follows from [7, Theorem 5.5] and [1, Theorem 2.6] that there exists a unique weak solution to (1.1). That is:

**Lemma 2.2.** For any  $f \in W^{-\alpha,p'}(\Omega)$ , the dual space of  $W^{\alpha,p}(\Omega)$ , there is a unique function  $u \in W^{\alpha,p}(\Omega)$  such that

 $L_{\alpha,p}u = f.$ 

The next lemma is a straightforward consequence of [1, Lemma 2.8, Proposition 2.10, Lemma 2.13], we collect here without their proofs.

# **Lemma 2.3.** Let $\mathfrak{u}, \nu \in \dot{W}^{\alpha, p}(\Omega)$ .

(a) If  $F \in Lip(\mathbb{R})$  with F(0) = 0, then  $F(u) \in \dot{W}^{\alpha,p}(\Omega)$ . Furthermore, if F is a convex function and differentiable almost everywhere, one has

$$L_{\alpha,p}F(\mathfrak{u}) \leq |F'(\mathfrak{u})|^{p-2}F'(\mathfrak{u})L_{\alpha,p}\mathfrak{u}$$
, a.e. in  $\Omega$ .

- (b)  $\langle L_{\alpha,p} \mathfrak{u}, \varphi \nu \rangle = 2 \int_{\Omega} \mathfrak{u} L_{\alpha,p} \nu$ .
- (c) For any  $\mathfrak{m} \ge 0$  and  $\xi \in \mathbb{R}$ , define the truncated functions

$$\mathfrak{F}_{\mathfrak{m}}(\xi) := \max\{-\mathfrak{m}, \min\{\mathfrak{m}, \xi\}\}, \text{ and } \mathfrak{H}_{\mathfrak{m}}(\xi) := \xi - \mathfrak{F}_{\mathfrak{m}}(\xi).$$

*Then*  $\mathcal{F}_{\mathfrak{m}}(\mathfrak{u}), \mathcal{H}_{\mathfrak{m}}(\mathfrak{u}) \in \dot{W}^{\alpha,p}(\Omega)$  and

$$|\mathfrak{F}_{\mathfrak{m}}(\mathfrak{u})||_{\dot{W}^{\alpha,p}(\Omega)}^{\mathfrak{p}} \leqslant \left\langle \mathsf{L}_{\alpha,p}\mathfrak{u},\mathfrak{F}_{\mathfrak{m}}(\mathfrak{u})\right\rangle, \quad and \quad \|\mathfrak{H}_{\mathfrak{m}}(\mathfrak{u})\|_{\dot{W}^{\alpha,p}(\Omega)}^{\mathfrak{p}} \leqslant \left\langle \mathsf{L}_{\alpha,p}\mathfrak{u},\mathfrak{H}_{\mathfrak{m}}(\mathfrak{u})\right\rangle.$$

The following fractional Morrey Sobolev inequality is essential to the proof of Theorem 1.2.

**Lemma 2.4** ([20, Theorem 4.1]). Let  $\alpha p > n$ . Then there is a constant  $C := C(n, \alpha, p, \Omega)$  such that

$$\|f\|_{L^{\infty}(\Omega)} \leq C \|f\|_{\dot{W}^{\alpha,p}(\Omega)}, \quad \forall f \in \dot{W}^{\alpha,p}(\Omega).$$

#### 3. Proofs of the main results

In this section, we give the proofs of Theorem 1.1 and Theorem 1.2 by the so-called Moser's method. Following [1], we can also give the proofs with a stampacchia's type result. We omit here for their similarity.

# 3.1. Proof of Theorem 1.1

Let us begin with the proof of (a). There is no loss of generality in assuming that  $|\Omega| = 1$ . We consider the following truncated function

$$F_{\Psi}(\xi) := \begin{cases} |\xi|^{\gamma} & \text{as } 0 \leq |\xi| < \Psi, \\ \gamma \Psi^{\gamma - 1}(\xi - \Psi) + \Psi^{\gamma} & \text{as } \xi \geq \Psi, \\ -\gamma \Psi^{\gamma - 1}(\xi + \Psi) + \Psi^{\gamma} & \text{as } \xi \leq -\Psi, \end{cases}$$
(3.1)

where  $\gamma \ge 1$  and  $\Psi > 0$  to be announced later. It is easy to check that  $F_{\Psi}$  satisfies Lemma 2.3 (a), and hence  $F_{\Psi}(\mathfrak{u}) \in \dot{W}^{\alpha,n/\alpha}(\Omega)$ . It follows from Lemma 2.1, (2.1) and Lemma 2.3 that

$$\begin{split} \|F_{\Psi}(u)\|_{L^{q}(\Omega)}^{n/\alpha} &\leqslant C \int_{\Omega} \int_{\Omega} \frac{|F_{\Psi}(u)(y) - F_{\Psi}(u)(x)|^{n/\alpha}}{|y - x|^{2n}} dy dx \\ &= C \langle L_{\alpha,n/\alpha} F_{\Psi}(u), F_{\Psi}(u) \rangle \\ &= 2C \int_{\Omega} F_{\Psi}(u)(x) \left[ L_{\alpha,n/\alpha} F_{\Psi}(u) \right] (x) dx \\ &\leqslant 2C \int_{\Omega} |F'_{\Psi}(u)(x)|^{n/\alpha - 1} F_{\Psi}(u)(x) L_{\alpha,n/\alpha} u(x) dx \\ &= 2C \int_{\Omega} |F'_{\Psi}(u)(x)|^{n/\alpha - 1} F_{\Psi}(u)(x) f(x) dx. \end{split}$$

Since

$$F_{\Psi}(\mathfrak{u}) \leqslant |\mathfrak{u}|^{\gamma}$$
, and  $|F'_{\Psi}(\mathfrak{u})| \leqslant \gamma |\mathfrak{u}|^{n/\alpha-1}$ ,

the Hölder inequality gives

$$\|F_{\Psi}(\mathfrak{u})\|_{L^{q}(\Omega)}^{n/\alpha} \leqslant 2C\gamma^{n/\alpha-1}\|g\|_{L^{s}(\Omega)} \left\||\mathfrak{u}|^{n/\alpha(\gamma-1)+1}\right\|_{L^{s'}(\Omega)}$$

Hence

$$\left\| |\mathfrak{u}|^{\gamma} \right\|_{L^{q}(\Omega)}^{n/\alpha} \leq C \gamma^{n/\alpha-1} \|g\|_{L^{s}(\Omega)} \left\| |\mathfrak{u}|^{n/\alpha(\gamma-1)+1} \right\|_{L^{s'}(\Omega)}$$

by taking  $\Psi \to \infty$ . Finally,

$$\|\mathbf{u}\|_{L^{\gamma \mathfrak{q}}(\Omega)} \leqslant C\left(\gamma^{n/\alpha-1} \|g\|_{L^{s}(\Omega)}\right)^{\frac{\alpha}{\gamma \mathfrak{n}}} \left(\int_{\Omega} |\mathbf{u}(\mathbf{x})|^{(n/\alpha(\gamma-1)+1)s'} d\mathbf{x}\right) \frac{\alpha}{\mathfrak{n}\gamma s'}.$$
(3.2)

Applying Young's inequality to  $p_1 = \frac{n\gamma}{n(\gamma-1)+\alpha}$  and  $p_2 = \frac{n\gamma}{n-\alpha}$ , (3.2) can be rewritten as

$$\|\mathbf{u}\|_{L^{\gamma \mathfrak{q}}(\Omega)} \leq C\left(\gamma^{n/\alpha-1} \|\mathbf{g}\|_{L^{s}(\Omega)}\right)^{\frac{\alpha}{\gamma n}} \left(1 + \int_{\Omega} |\mathbf{u}(x)|^{n\gamma s'/\alpha} dx\right) \frac{\alpha}{n\gamma s'}.$$
(3.3)

Denote by

$$\mathbf{r}_{k} = \mathbf{m}^{k} := \left(\frac{q\alpha}{\mathbf{n}s'}\right)^{k}, \quad \mathbf{I}_{k} = \left(\int_{\Omega} |\mathbf{u}|^{\mathbf{r}_{k}q} d\mathbf{x}\right)^{\frac{1}{\mathbf{r}_{k}q}}, \text{ and } \mathbf{J}_{k} = \left(c\mathbf{r}_{k}^{\mathbf{n}/\alpha-1} \|\mathbf{f}\|_{L^{s}(\Omega)}\right)^{\frac{\alpha}{\mathbf{r}_{k}n}}.$$

It is obvious that m > 1. We conclude from the fact  $r_{k+1}\frac{ns'}{\alpha} = r_k q$  and (3.3) that

$$I_{k+1} \leqslant J_{k+1} (1 + I_k^{r_k q})^{\frac{1}{r_k q}}.$$

Therefore, up to a re-normalization to obtain that  $I_0 = 1$  and  $I_k \ge 1$ , one has

$$ln^{I_{k+1}} \leqslant ln^{J_{k+1}} + \frac{1}{r_k q} \, ln^{(1+I_k^{r_k q})} \leqslant ln^{J_{k+1}} + \frac{1}{r_k q} + ln^{I_k} \, .$$

Hence,

$$\sum_{i=0}^{k+1} (ln^{I_{i+1}} - ln^{I_i}) \leqslant \sum_{i=0}^{k+1} \left( ln^{J_{i+1}} + \frac{1}{r_i q} \right),$$

which implies that

$$ln^{I_{k+1}} \leqslant \sum_{i=1}^{k+1} ln^{J_i} + \sum_{i=0}^k \frac{1}{r_i q} + ln^{I_0} \leqslant \sum_{i=1}^\infty ln^{J_i} + \sum_{i=0}^\infty \frac{1}{r_i q} = N < \infty.$$

We have completed our proof after the following observation

$$I_{k+1}\leqslant C:=e^N<\infty, \ \text{ and } \ \lim_{k\to\infty}I_{k+1}=\|u\|_{L^\infty(\Omega)}.$$

Next, we proceed the proof by showing (b). For any  $\Psi > 0$  and  $\gamma > 0$  will be fixed later, we define

$$\mathsf{G}_{\Psi}(\eta) := \begin{cases} e^{\gamma |\eta|} - 1 & \text{ as } 0 \leqslant |\eta| < \Psi, \\ \gamma e^{\gamma \Psi} (\eta - \Psi) + e^{\gamma \Psi} - 1 & \text{ as } \eta \geqslant \Psi, \\ -\gamma e^{\gamma \Psi} (\eta + \Psi) + e^{\gamma \Psi} - 1 & \text{ as } \eta \leqslant -\Psi. \end{cases}$$

It is easily seen that  $G_{\Psi}(u)$  satisfies Lemma 2.3 (a). A further use of Lemma 2.1 and Lemma 2.3, one has

$$\begin{split} \|G_{\Psi}(u)\|_{L^{q}(\Omega)}^{n/\alpha} &\leqslant C\langle L_{\alpha,n/\alpha}G_{\Psi}(u),G_{\Psi}(u)\rangle \\ &= 2C\int_{\Omega}G_{\Psi}(u)(x)\left[L_{\alpha,n/\alpha}G_{\Psi}(u)\right](x)dx \\ &\leqslant 2C\int_{\Omega}|G_{\Psi}'(u)(x)|^{n/\alpha-1}G_{\Psi}(u)(x)L_{\alpha,n/\alpha}u(x)dx \\ &= 2C\int_{\Omega_{1}^{c}}|G_{\Psi}'(u)(x)|^{n/\alpha-1}G_{\Psi}(u)(x)g(x)dx \\ &\quad +\int_{\Omega_{1}}|G_{\Psi}'(u)(x)|^{n/\alpha-1}G_{\Psi}(u)(x)g(x)dx \\ &:= 2C(K_{1}+K_{2}), \end{split}$$

where  $\Omega_1 := \{ x \in \Omega : |u(x)| \ge \Psi \}.$ 

Set

$$\Omega_{11}^{c}:=\{x\in\Omega_{1}^{c}:G_{\Psi}(\mathfrak{u})(x)\geqslant1\}, \text{ and } \Omega_{12}^{c}:=\{x\in\Omega_{1}^{c}:G_{\Psi}(\mathfrak{u})(x)<1\}.$$

We see at once that

$$(G_{\Psi}(\mathfrak{u})(x)+1)^{n/\alpha-1} \leqslant 2^{n/\alpha-1} \left( (G_{\Psi}(\mathfrak{u})(x))^{n/\alpha-1}+1 \right) \text{ as } x \in \Omega_{11}^c,$$

and

$$(G_{\Psi}(\mathfrak{u})(x)+1)^{n/\alpha-1}\leqslant 2^{n/\alpha-1}\leqslant 2^{n/\alpha-1}\big((G_{\Psi}(\mathfrak{u})(x))^{n/\alpha-1}+1\big) \text{ as } x\in \Omega_{12}^c.$$

Since

$$|\mathsf{G}'_{\Psi}(\mathfrak{u})(\mathfrak{x})| = \gamma(\mathsf{G}_{\Psi}(\mathfrak{u})(\mathfrak{x}) + 1) \text{ as } \mathfrak{x} \in \Omega_1^c$$

the Hölder inequality shows that

$$\begin{split} \mathsf{K}_1 &\leqslant 2^{n/\alpha - 1} \gamma^{n/\alpha - 1} \int_{\Omega_1^c} \left[ (\mathsf{G}_{\Psi}(\mathfrak{u})(x))^{n/\alpha} g(x) + (\mathsf{G}_{\Psi}(\mathfrak{u})(x)) g(x) \right] dx \\ &\leqslant (2\gamma)^{n/\alpha - 1} \|g\|_{\mathsf{L}^s(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^q(\Omega)}^{n/\alpha} + (2\gamma)^{n/\alpha - 1} \int_{\Omega} \mathsf{G}_{\Psi}(\mathfrak{u})(x) g(x) dx \\ &\coloneqq \mathsf{K}_{11} + \mathsf{K}_{12}, \end{split}$$

and hence

$$\mathsf{K}_{12} \leqslant (2\gamma)^{n/\alpha-1} \|g\|_{\mathsf{L}^{\mathfrak{s}}(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathfrak{s}'}(\Omega)} \leqslant (2\gamma)^{n/\alpha-1} \|g\|_{\mathsf{L}^{\mathfrak{s}}(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathfrak{q}}(\Omega)} |\Omega|^{\frac{n-\alpha}{\alpha \mathfrak{q}}}.$$

By the Minkowski inequality,  $K_{12}\xspace$  can be further estimated as

$$\mathsf{K}_{12} \leqslant (2\gamma)^{n/\alpha-1} \left( \frac{\alpha}{n} \|g\|_{\mathsf{L}^{\mathfrak{s}}(\Omega)} \|G_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathfrak{q}}(\Omega)}^{n/\alpha} |\Omega|^{\frac{n(n-\alpha)}{\alpha^{2}\mathfrak{q}}} + \frac{n-\alpha}{n} \|g\|_{\mathsf{L}^{\mathfrak{s}}(\Omega)} \right).$$

Therefore,

$$\begin{split} \mathsf{K}_1 &\leqslant 2^{n/\alpha} \gamma^{n/\alpha-1} \|g\|_{L^s(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{L^q(\Omega)}^{n/\alpha} + \frac{\alpha}{n} \|g\|_{L^s(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{L^q(\Omega)}^{n/\alpha} |\Omega|^{\frac{n(n-\alpha)}{\alpha^2 q}} \\ &+ \frac{n-\alpha}{n} \|g\|_{L^s(\Omega)} \\ &\leqslant 2^{n/\alpha} \gamma^{n/\alpha-1} \|g\|_{L^s(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{L^q(\Omega)}^{n/\alpha} \left(1 + \frac{\alpha}{n} |\Omega|^{\frac{n(n-\alpha)}{\alpha^2 q}}\right) \\ &+ 2^{n/\alpha} \gamma^{n/\alpha-1} \frac{n-\alpha}{n} \|g\|_{L^s(\Omega)}. \end{split}$$

For the term  $K_2$ , we first note that

$$\min_{\Omega_1} \mathsf{G}_{\Psi}(\mathfrak{u}) = \mathsf{G}_{\Psi}(\Psi) = \mathsf{G}_{\Psi}(-\Psi) = e^{\gamma \Psi} - 1.$$

Choosing  $\gamma \Psi > 1$ , Hölder's inequality gives

$$\begin{split} \mathsf{K}_{2} &\leqslant \frac{(\gamma e^{\gamma \Psi})^{n/\alpha - 1}}{\mathsf{G}_{\Psi}(\Psi)^{n/\alpha - 1}} \int_{\Omega_{1}} \mathsf{G}_{\Psi}(\mathfrak{u}(x))^{n/\alpha} \mathfrak{g}(x) dx \\ &\leqslant \gamma^{n/\alpha - 1} \left( \frac{e^{\gamma \Psi}}{e^{\gamma \Psi} - 1} \right)^{n/\alpha - 1} \|\mathfrak{g}\|_{\mathsf{L}^{s}(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathfrak{q}}(\Omega)}^{n/\alpha} \\ &\leqslant (2\gamma)^{n/\alpha - 1} \|\mathfrak{g}\|_{\mathsf{L}^{s}(\Omega)} \|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathfrak{q}}(\Omega)}^{n/\alpha}. \end{split}$$

On account of the above estimates, we have

$$\begin{split} \|G_{\Psi}(\mathfrak{u})\|_{L^{q}(\Omega)}^{n/\alpha} &\leq 2^{n/\alpha}(\gamma)^{n/\alpha-1} \|g\|_{L^{s}(\Omega)} \|G_{\Psi}(\mathfrak{u})\|_{L^{q}(\Omega)}^{n/\alpha} \left(2 + \frac{\alpha}{n} |\Omega|^{\frac{n(n-\alpha)}{\alpha^{2}q}}\right) \\ &+ \frac{n-\alpha}{n} 2^{n/\alpha}(\gamma)^{n/\alpha-1} \|g\|_{L^{s}(\Omega)}. \end{split}$$

Therefore,

$$\|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathsf{q}}(\Omega)}^{\mathfrak{n}/\alpha} \leqslant C^{\frac{\mathfrak{n}}{\alpha}}(\mathfrak{n},\alpha,\mathfrak{q},\|\mathfrak{g}\|_{\mathsf{L}^{\mathfrak{s}}(\Omega)},\Omega),$$

by choosing  $\gamma$  small enough. This forces

$$|\mathsf{G}_{\Psi}(\mathfrak{u})\|_{\mathsf{L}^{\mathfrak{q}}(\Omega)} = \|e^{\gamma|\mathfrak{u}|} - 1\|_{\mathsf{L}^{\mathfrak{q}}(\Omega)} \leqslant C < \infty,$$

after taking  $\Psi \to \infty$ . Thus, we get our desired result by taking  $\beta = \gamma q > 0$ .

Finally, we give the proof of (c). For  $\Psi > 0$  big enough to be determined later, we consider the function (3.1) with  $\gamma = \frac{s^*}{q}$ . We deduce from the fact  $\frac{q}{q-1} \leq s < \frac{\alpha q}{\alpha q-n}$  that  $\gamma \geq 1$ . By the similar analysis as that in the proof of (a), we have

$$\|F_{\Psi}(\mathfrak{u})\|_{L^{q}(\Omega)}^{n/\alpha} \leq 2C\gamma^{\frac{n}{\alpha}-1}|g\|_{L^{s}(\Omega)} \left(\int_{\Omega} |\mathfrak{u}(x)|^{s'((\gamma-1)\frac{n}{\alpha}+1)} dx\right)^{\frac{1}{s'}}.$$

Hence,

$$\left(\int_{\Omega} |u(x)|^{\gamma q} dx\right)^{\frac{n}{q-\alpha}} \leq 2C\gamma^{\frac{n}{\alpha}-1} \|g\|_{L^{s}(\Omega)} \left(\int_{\Omega} |u(x)|^{s'((\gamma-1)\frac{n}{\alpha}+1)} dx\right)^{\frac{1}{s'}},$$

by taking  $\Psi \to \infty$ . It follows from the fact

$$\gamma q = s'((\gamma - 1)\frac{n}{\alpha} + 1) = s^*, \quad \text{and} \quad \frac{n}{\alpha q} - \frac{1}{s'} = \frac{\alpha q - s(\alpha q - n)}{\alpha q s} > 0,$$
$$\|u\|^{\frac{n-\alpha}{\alpha}} \leq 2C \gamma^{n/\alpha - 1} \|q\|_{L^{\infty}(\Omega)}.$$

that

$$\|\mathbf{u}\|_{L^{s^*}(\Omega)}^{\frac{n-\alpha}{\alpha}} \leq 2C\gamma^{n/\alpha-1} \|\mathbf{g}\|_{L^s(\Omega)},$$

which gives

$$\|\mathfrak{u}\|_{L^{s^*}(\Omega)} \leqslant C_3 \|g\|_{L^s(\Omega)}^{\frac{\alpha}{n-\alpha}}.$$

# 3.2. Proof of Theorem 1.2

Using the similar arguments as that of Section 3.1, we can prove Theorem 1.2. In fact, for the function defined in (3.1) with  $\gamma = 1$  and  $\Psi > 0$ , by Lemma 2.3 and Lemma 2.4, we have

$$\begin{split} \|F_{\Psi}(u)\|_{L^{\infty}(\Omega)} &\leqslant 2\tilde{C} \left( \int_{\Omega} |F'_{\Psi}(u(x))|^{p-1} F_{\Psi}(u(x))g(x)dx \right)^{1/p} \\ &\leqslant 2\tilde{C} \|u\|_{L^{\infty}(\Omega)}^{1/p} \|g\|_{L^{1}(\Omega)}^{1/p}. \end{split}$$

Taking  $\Psi \to \infty$ , one has

$$\|\mathbf{u}\|_{L^{\infty}(\Omega)} \leq 2\tilde{C} \|\mathbf{u}\|_{L^{\infty}(\Omega)}^{1/p} \|g\|_{L^{1}(\Omega)}^{1/p}$$

Hence

$$\|\mathfrak{u}\|_{L^{\infty}(\Omega)} \leqslant 2\tilde{C} \|g\|_{L^{1}(\Omega)}^{p'/p}$$

Theorem 1.2 has been proved by taking  $C = 2\tilde{C} \|g\|_{L^1(\Omega)}^{p'/p}$ .

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